

generators t^b of a representation of the gauge group. The generators obey the commutation relations

$$[t^a, t^b] = if_{abc}t^c \quad (10.508)$$

in which the f_{abc} are the structure constants of the gauge group. Show that under a gauge transformation (10.474)

$$A'_i = UA_iU^{-1} - (\partial_i U)U^{-1} \quad (10.509)$$

by the unitary matrix $U = \exp(-ig\lambda^a t^a)$ in which λ^a is infinitesimal, the gauge-field matrix A_i transforms as

$$-igA_i'^a t^a = -igA_i^a t^a - ig^2 f_{abc} \lambda^a A_i^b t^c + ig\partial_i \lambda^a t^a. \quad (10.510)$$

Show further that the gauge field transforms as

$$A_i'^a = A_i^a - \partial_i \lambda^a - gf_{abc} A_i^b \lambda^c. \quad (10.511)$$

- 10.43 Show that if the vectors $e_a(x)$ are orthonormal, then $e^{a\dagger} \cdot e_{c,i} = -e_{i,i}^{a\dagger} \cdot e_c$.
 10.44 Use the identity of exercise 10.43 to derive the formula (10.489) for the nonabelian Faraday tensor.

- 10.45 Write Dirac's action density in the explicitly hermitian form $L_D = -\frac{1}{2}\bar{\psi}\gamma^i\partial_i\psi - \frac{1}{2}[\bar{\psi}\gamma^i\partial_i\psi]^\dagger$ in which the field ψ has the invariant form $\psi = e_a\psi_a$ and $\bar{\psi} = i\psi^\dagger\gamma^0$. Use the identity $[\bar{\psi}_a\gamma^i\psi_b]^\dagger = -\bar{\psi}_b\gamma^i\psi_a$ to show that the gauge-field matrix A_i defined as the coefficient of $\bar{\psi}_a\gamma^i\psi_b$ as in $\bar{\psi}_a\gamma^i(\partial_i + iA_{iab})\psi_b$ is hermitian $A_{iab}^* = A_{iba}$.