

Goldstone bosons & the Higgs mechanism

First, Goldstone bosons:

$$\mathcal{L} = \sum_{i=1}^2 -\frac{1}{2} \partial_\mu \phi_i \partial^\mu \phi_i - \lambda \left(\sum \phi_i^2 - \phi_0^2 \right)^2$$

Must have

$$\langle \sum \phi_i^2 \rangle_0 = \phi_0^2$$

Suppose $\phi_i = \phi_0 + \tilde{\phi}_i$, that is

$$\langle \phi_1 \rangle_0 = \phi_0 \quad \text{and} \quad \langle \phi_2 \rangle_0 = 0.$$

Then

$$\begin{aligned} \phi_1^2 + \phi_2^2 - \phi_0^2 &= (\phi_0 + \tilde{\phi}_1)^2 + \phi_2^2 - \phi_0^2 \\ &= 2\phi_0 \tilde{\phi}_1 + \tilde{\phi}_1^2 + \phi_2^2, \quad \text{so} \end{aligned}$$

$$\begin{aligned} \mathcal{L} &= -\frac{1}{2} \partial_\mu \tilde{\phi}_1 \partial^\mu \tilde{\phi}_1 - \frac{1}{2} \partial_\mu \phi_2 \partial^\mu \phi_2 \\ &\quad - \lambda \left(2\phi_0 \tilde{\phi}_1 + \tilde{\phi}_1^2 + \phi_2^2 \right)^2 \end{aligned}$$

$$\begin{aligned} &= \mathcal{L}_0 - 4\lambda \phi_0^2 \tilde{\phi}_1^2 - 4\lambda \phi_0 \tilde{\phi}_1 (\tilde{\phi}_1^2 + \phi_2^2) \\ &\quad - \lambda (\tilde{\phi}_1^2 + \phi_2^2)^2 \end{aligned}$$

The usual mass term is $-\frac{m^2}{2} \phi^2$,
so

$$-\frac{m^2}{2} = -4\lambda \phi_0^2 \quad m^2 = 8\lambda \phi_0^2$$

So in this theory, ϕ_1 describes a boson of mass $2\sqrt{2}\lambda\phi_0$ and ϕ_2 a boson of mass zero. The massless boson is called a Goldstone boson.

Now, the Higgs mechanism:

$$\mathcal{L}_2 = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} (D_\mu \phi)_i (D^\mu \phi)_i - \lambda (\sum \phi_i^2 - \phi_0^2)^2,$$

where

$$D_\mu \phi = \partial_\mu \phi + e A_\mu t$$

and t is the matrix

$$t = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \text{ with } t^2 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = -I$$

which generates $O(2) = U(1)$,

$$\phi' = e^{\theta t} \phi = \left[\sum_{n=0}^{\infty} \frac{\theta^{2n}}{(2n)!} + t \sum_{n=0}^{\infty} \frac{\theta^{2n+1}}{(2n+1)!} \right] \phi$$

$$\phi' = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \phi$$

Now the Higgs potential $\propto (\phi_1^2 - \phi_2^2)^2$ still gives mass $m = g + \phi_0^2$ to $\tilde{\phi}_1$'s

boson, but also the $(D_\mu \phi)^2$ term is a mass term for the photon:

$$\begin{aligned} -\frac{1}{2} (D_\mu \phi)^\dagger D^\mu \phi &= -\frac{1}{2} \phi^\dagger (\overleftarrow{\partial}_\mu + e A_\mu \tau^T) (\partial^\mu + e A^\mu \tau) \phi \\ &= -\frac{1}{2} (\phi_0 + \tilde{\phi}_1, \phi_2) (\overleftarrow{\partial}_\mu + e A_\mu \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}) (\partial^\mu + e A^\mu \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}) \begin{pmatrix} \phi_0 + \tilde{\phi}_1 \\ \phi_2 \end{pmatrix} \\ &= -\frac{1}{2} (\partial_\mu \tilde{\phi}_1 + e A_\mu \phi_2, \partial_\mu \phi_2 - e A_\mu (\phi_0 + \tilde{\phi}_1)) \begin{pmatrix} \partial^\mu \tilde{\phi}_1 + e A^\mu \phi_2 \\ \partial^\mu \phi_2 - e A^\mu (\phi_0 + \tilde{\phi}_1) \end{pmatrix} \\ &= -\frac{1}{2} (\partial_\mu \tilde{\phi}_1 + e A_\mu \phi_2) (\partial^\mu \tilde{\phi}_1 + e A^\mu \phi_2) \\ &\quad - \frac{1}{2} (\partial_\mu \phi_2 - e A_\mu (\phi_0 + \tilde{\phi}_1)) (\partial^\mu \phi_2 - e A^\mu (\phi_0 + \tilde{\phi}_1)) \end{aligned}$$

in which $-\frac{1}{2} e^2 \phi_0^2 A_\mu A^\mu$ is a mass term

for A_μ . The full mass term is

$$-\frac{1}{2} e^2 \phi_0^2 W_\mu W^\mu \quad \text{where} \quad W_\mu = A_\mu - \frac{1}{e \phi_0} \partial_\mu \phi_2.$$

So the "longitudinal" component of

$$W_\mu \text{ is } -\frac{1}{e\phi_0} \partial_\mu \phi_2.$$

Note that

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu = \partial_\mu W_\nu - \partial_\nu W_\mu$$

because $[\partial_\mu, \partial_\nu] \phi = 0.$

So instead of a theory with 2 spinless bosons, one massive, and one massless photon, we have one massive spinless boson, "the Higgs," and one massive spin-1 boson, "the W." This is the "Higgs mechanism." In the standard model the gauge group is $SU_2^L \otimes U_1$,

in which the L means that only the "left-handed" fermions are mixed by SU_2 .