

$$\vec{J} = \vec{L} + \vec{S}$$

An electron in a central potential has orbital angular momentum \vec{L} and spin \vec{S} . Its angular momentum states are

$$|l, m\rangle \otimes |l, \pm\rangle.$$

These are eigenstates of \vec{L}^2 , L_z , \vec{S}^2 , and S_z .

The electron's total angular momentum is

$$\vec{J} = \vec{L} + \vec{S}.$$

Its eigenstates of \vec{J}^2 , J_z , \vec{L}^2 , and \vec{S}^2 are

$$|J, M\rangle = |J, M, l, \pm\rangle.$$

with eigenvalues

$$\vec{J}^2 |J, M\rangle = \hbar^2 J(J+1) |J, M\rangle$$

$$J_z |J, M\rangle = \hbar M |J, M\rangle$$

$$\vec{L}^2 |J, M, l, \pm\rangle = \hbar^2 l(l+1) |J, M, l, \pm\rangle$$

$$\begin{aligned} \vec{S}^2 |J, M, l, \pm\rangle &= \hbar^2 \frac{1}{2} \left(\frac{3}{2} \right) |J, M, l, \pm\rangle \\ &= \frac{3}{4} \hbar^2 |J, M, l, \pm\rangle. \end{aligned}$$

The state of highest J and M is

$$|JM\rangle = |l+\frac{1}{2}, l+\frac{1}{2}\rangle = |l, l\rangle |l, +\rangle.$$

Recall $J_- = J_1 - iJ_2$ and

$$\begin{aligned} J_- |j, m\rangle &= \hbar \sqrt{j(j+1) - m(m-1)} |j, m-1\rangle \\ &= \hbar \sqrt{(j+m)(j+1-m)} |j, m-1\rangle. \end{aligned}$$

We use J_- to find $|l+\frac{1}{2}, l-\frac{1}{2}\rangle$:

$$\begin{aligned}
 J_- |l+\frac{1}{2}, l+\frac{1}{2}\rangle &= \hbar \sqrt{(2l+1)} |l+\frac{1}{2}, l-\frac{1}{2}\rangle \\
 &= (L_- + S_-) |l, l\rangle |\frac{1}{2}, +\rangle \\
 &= (L_-(l, l\rangle) |\frac{1}{2}, +\rangle + |l, l\rangle S_- |\frac{1}{2}, +\rangle) \\
 &= \hbar \sqrt{2l} (|l, l-1\rangle |\frac{1}{2}, +\rangle + |l, l\rangle \hbar |\frac{1}{2}, -\rangle).
 \end{aligned}$$

Thus

$$|l+\frac{1}{2}, l-\frac{1}{2}\rangle = \frac{1}{\sqrt{2l+1}} (\sqrt{2l} |l, l-1\rangle |\frac{1}{2}, +\rangle + |l, l\rangle |\frac{1}{2}, -\rangle).$$

Now we use J_- again to find
 $|l+\frac{1}{2}, l-\frac{3}{2}\rangle$:

$$\begin{aligned}
 J_- |l+\frac{1}{2}, l-\frac{1}{2}\rangle &= \hbar \sqrt{2l \cdot 2} |l+\frac{1}{2}, l-\frac{3}{2}\rangle \\
 &= \frac{(L_- + S_-)}{\sqrt{2l+1}} (\sqrt{2l} |l, l-1\rangle |\frac{1}{2}, +\rangle + |l, l\rangle |\frac{1}{2}, -\rangle) \\
 &= \frac{1}{\sqrt{2l+1}} \left(\sqrt{2l} L_-(l, l-1\rangle |\frac{1}{2}, +\rangle + \sqrt{2l} |l, l-1\rangle S_- |\frac{1}{2}, +\rangle \right. \\
 &\quad \left. + L_- |l, l\rangle |\frac{1}{2}, -\rangle + |l, l\rangle S_- |\frac{1}{2}, -\rangle \right).
 \end{aligned}$$

But $S_- |\frac{1}{2}, -\rangle = 0$, and the other terms give

$$\langle J=1, l+\frac{1}{2}, l-\frac{1}{2} \rangle = \frac{1}{\sqrt{2l+1}} \left(\sqrt{2l} + \sqrt{(2l-1) \cdot 2} |l, l-2\rangle |1\frac{1}{2}, +\rangle + \sqrt{2l} |l, l-1\rangle |1\frac{1}{2}, -\rangle + \sqrt{2l} |l, l-1\rangle |1\frac{1}{2}, -\rangle \right).$$

Thus

$$\begin{aligned} |l+\frac{1}{2}, l-\frac{3}{2}\rangle &= \frac{1}{\sqrt{4l}} \frac{1}{\sqrt{2l+1}} \left(\sqrt{2l} \sqrt{2(2l-1)} |l, l-2\rangle |1\frac{1}{2}, +\rangle + 2\sqrt{2l} |l, l-1\rangle |1\frac{1}{2}, -\rangle \right. \\ &\quad \left. = \frac{1}{\sqrt{2l+1}} \left(\sqrt{2l-1} |l, l-2\rangle |1\frac{1}{2}, +\rangle + \sqrt{2} |l, l-1\rangle |1\frac{1}{2}, -\rangle \right) \right). \end{aligned}$$

The other $|l+\frac{1}{2}, M\rangle$ states are

$$\begin{aligned} |l+\frac{1}{2}, M\rangle &= \frac{1}{\sqrt{2l+1}} \left(\sqrt{l+M+\frac{1}{2}} |l, M-\frac{1}{2}\rangle |1\frac{1}{2}, +\rangle \right. \\ &\quad \left. + \sqrt{l-M+\frac{1}{2}} |l, M+\frac{1}{2}\rangle |1\frac{1}{2}, -\rangle \right) \end{aligned}$$

where

$$M = l+\frac{1}{2}, l-\frac{1}{2}, l-\frac{3}{2}, \dots -l+\frac{1}{2}, -(l+\frac{1}{2}).$$

These are all the states with $J=l+\frac{1}{2}$.

The other states have $J=l-\frac{1}{2}$.

To find them, we know that the state

$|l-\frac{1}{2}, l-\frac{1}{2}\rangle$ must be the linear combination of the states $|l, l\rangle |l\frac{1}{2}, -\rangle$ and $|l, l-1\rangle |l\frac{1}{2}, +\rangle$ that is orthogonal to the state

$$|l+\frac{1}{2}, l-\frac{1}{2}\rangle = \overline{\frac{1}{\sqrt{2l+1}}} \left(\sqrt{2l} |l, l-1\rangle |l\frac{1}{2}, +\rangle + |l, l\rangle |l\frac{1}{2}, -\rangle \right).$$

By convention, we choose

$$|l-\frac{1}{2}, l-\frac{1}{2}\rangle = \overline{\frac{1}{\sqrt{2l+1}}} \left(\sqrt{2l} |l, l\rangle |l\frac{1}{2}, -\rangle - |l, l-1\rangle |l\frac{1}{2}, +\rangle \right).$$

Similarly,

$$|l-\frac{1}{2}, l-\frac{3}{2}\rangle = \overline{\frac{1}{\sqrt{2l+1}}} \left(\sqrt{2l-1} |l, l-1\rangle |l\frac{1}{2}, -\rangle - \sqrt{2} |l, l-2\rangle |l\frac{1}{2}, +\rangle \right).$$

More generally,

$$|l-\frac{1}{2}, M\rangle = \overline{\frac{1}{\sqrt{2l+1}}} \left(\sqrt{l+M+\frac{1}{2}} |l, M+\frac{1}{2}\rangle |l\frac{1}{2}, -\rangle - \sqrt{l-M+\frac{1}{2}} |l, M-\frac{1}{2}\rangle |l\frac{1}{2}, +\rangle \right).$$

for

$$M = l-\frac{1}{2}, l-\frac{3}{2}, \dots, -l+\frac{3}{2}, -(l-\frac{1}{2}).$$