

Orbital Angular Momentum in Spherical Coordinates

A change $d\vec{r}$ in the position is

$$\begin{aligned}d\vec{r} &= \hat{x} dx + \hat{y} dy + \hat{z} dz \\ &= \hat{r} dr + r\hat{\theta} d\theta + r\sin\theta\hat{\phi} d\phi \quad (1)\end{aligned}$$

in euclidean and spherical coordinates.

In terms of \hat{x} , \hat{y} , & \hat{z} , the vectors \hat{r} , $\hat{\theta}$, & $\hat{\phi}$ are

$$\hat{r} = \hat{x} \sin\theta \cos\phi + \hat{y} \sin\theta \sin\phi + \hat{z} \cos\theta$$

$$\hat{\theta} = \hat{x} \cos\theta \cos\phi + \hat{y} \cos\theta \sin\phi - \hat{z} \sin\theta$$

$$\hat{\phi} = -\hat{x} \sin\phi + \hat{y} \cos\phi. \quad (2)$$

The inverse relations are

$$\hat{x} = \hat{x} \cdot \hat{r} \hat{r} + \hat{x} \cdot \hat{\theta} \hat{\theta} + \hat{x} \cdot \hat{\phi} \hat{\phi}$$

$$\hat{x} = \sin\theta \cos\phi \hat{r} + \cos\theta \cos\phi \hat{\theta} - \sin\phi \hat{\phi}$$

$$\hat{y} = \sin\theta \sin\phi \hat{r} + \cos\theta \sin\phi \hat{\theta} + \cos\phi \hat{\phi}$$

$$\hat{z} = \cos\theta \hat{r} - \sin\theta \hat{\theta}, \quad (3)$$

The gradient $\vec{\nabla} f$ is defined by

$$df = \vec{dr} \cdot \vec{\nabla} f = dr \frac{\partial f}{\partial r} + d\theta \frac{\partial f}{\partial \theta} + d\phi \frac{\partial f}{\partial \phi} \quad (4)$$

and is in view of (1)

$$\vec{\nabla} f = \hat{r} \frac{\partial f}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial f}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial f}{\partial \phi} \quad (5)$$

So

$$\vec{\nabla} = \hat{r} \frac{\partial}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial}{\partial \phi} \quad (6)$$

and the orbital angular momentum

$$\vec{L} = \vec{r} \times \vec{p} \quad (7)$$

is represented as

$$\langle \vec{r} | \vec{L} | \psi \rangle = \vec{r} \times \frac{\hbar}{i} \vec{\nabla} \langle \vec{r} | \psi \rangle \quad (8)$$

$$= \frac{\hbar}{i} \hat{r} \times \left(\hat{r} \frac{\partial}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \langle \vec{r} | \psi \rangle$$

$$= \frac{\hbar}{i} \left(\hat{\phi} \frac{\partial}{\partial \theta} - \frac{\hat{\theta}}{\sin \theta} \frac{\partial}{\partial \phi} \right) \langle \vec{r} | \psi \rangle \quad (9)$$

The rectangular components of \vec{L} are represented by

$$\begin{aligned} \langle \vec{r} | L_z | \psi \rangle &= \langle \vec{r} | \hat{z} \cdot \vec{L} | \psi \rangle \\ &= \hat{z} \cdot \frac{\hbar}{i} \left(\hat{\phi} \partial_\theta - \frac{\hat{\theta}}{\sin \theta} \partial_\phi \right) \langle \vec{r} | \psi \rangle \\ &= \frac{\hbar}{i} \partial_\phi \langle \vec{r} | \psi \rangle \quad (10) \end{aligned}$$

and, less simply, by

$$\begin{aligned} \langle \vec{r} | L_x | \psi \rangle &= \langle \vec{r} | \hat{x} \cdot \vec{L} | \psi \rangle \\ &= \hat{x} \cdot \frac{\hbar}{i} \left(\hat{\phi} \partial_\theta - \frac{\hat{\theta}}{\sin \theta} \partial_\phi \right) \langle \vec{r} | \psi \rangle \\ &= \frac{\hbar}{i} \left(\cos \theta \cos \phi \hat{\theta} - \sin \phi \hat{\phi} \right) \cdot \left(\hat{\phi} \partial_\theta - \frac{\hat{\theta}}{\sin \theta} \partial_\phi \right) \langle \vec{r} | \psi \rangle \\ &= \frac{\hbar}{i} \left(-\sin \phi \partial_\theta - \cot \theta \cos \phi \partial_\phi \right) \langle \vec{r} | \psi \rangle \quad (11) \end{aligned}$$

and

$$\begin{aligned} \langle \vec{r} | L_y | \psi \rangle &= \langle \vec{r} | \hat{y} \cdot \vec{L} | \psi \rangle = \hat{y} \cdot \frac{\hbar}{i} \left(\hat{\phi} \partial_\theta - \frac{\hat{\theta}}{\sin \theta} \partial_\phi \right) \langle \vec{r} | \psi \rangle \\ &= \frac{\hbar}{i} \left(\cos \theta \sin \phi \hat{\theta} + \cos \phi \hat{\phi} \right) \cdot \left(\hat{\phi} \partial_\theta - \frac{\hat{\theta}}{\sin \theta} \partial_\phi \right) \langle \vec{r} | \psi \rangle \end{aligned}$$

that is,

$$\langle \vec{r} | L_y | \psi \rangle = \frac{\hbar}{i} \left(\cos \theta \partial_\theta - \cot \theta \sin \theta \partial_\phi \right) \langle \vec{r} | \psi \rangle. \quad (12)$$

It follows from (11) & (12) that the raising and lowering operators L_\pm are

$$\begin{aligned} \langle \vec{r} | L_\pm | \psi \rangle &= \langle \vec{r} | L_x \pm i L_y | \psi \rangle \\ &= \frac{\hbar}{i} \left[(-\sin \theta \pm i \cos \theta) \partial_\theta - \cot \theta (-\cos \theta \pm i \sin \theta) \partial_\phi \right. \\ &\quad \left. - \cot \theta (\cos \theta \pm i \sin \theta) \partial_\phi \right] \langle \vec{r} | \psi \rangle \\ &= \frac{\hbar}{i} e^{\pm i \phi} \left(\pm i \partial_\theta - \cot \theta \partial_\phi \right) \langle \vec{r} | \psi \rangle. \quad (13) \end{aligned}$$

Since $L^2 = L_z^2 + \frac{1}{2}(L_+ L_- + L_- L_+)$ (14)

one can show that

$$\langle \vec{r} | L^2 | \psi \rangle = -\hbar^2 \left[\frac{1}{\sin^2 \theta} \partial_\phi^2 + \frac{1}{\sin \theta} \partial_\theta (\sin \theta \partial_\theta) \right] \langle \vec{r} | \psi \rangle \quad (15)$$

by using (10) and (13).

In spherical coordinates, the Laplacian

$$\Delta = \vec{\nabla} \cdot \vec{\nabla} = \sum_{i=1}^3 \frac{\partial^2}{\partial x_i^2} \quad (16)$$

is

$$\Delta = \frac{1}{r^2} \left[\partial_r (r^2 \partial_r) + \frac{\partial_\theta (\sin \theta \partial_\theta)}{\sin \theta} + \frac{\partial_\phi^2}{\sin^2 \theta} \right] \quad (17)$$

Thus

$$\begin{aligned} \langle \vec{r} | \vec{p}^2 | \psi \rangle &= -\hbar^2 \nabla \cdot \nabla \langle \vec{r} | \psi \rangle = -\hbar^2 \Delta \langle \vec{r} | \psi \rangle \\ &= -\frac{\hbar^2}{r^2} \left[\partial_r (r^2 \partial_r) + \frac{\partial_\theta (\sin \theta \partial_\theta)}{\sin \theta} + \frac{\partial_\phi^2}{\sin^2 \theta} \right] \langle \vec{r} | \psi \rangle \\ &= -\frac{\hbar^2}{r^2} \partial_r (r^2 \partial_r) \langle \vec{r} | \psi \rangle + \frac{1}{r^2} \langle \vec{r} | \vec{L}^2 | \psi \rangle. \quad (18) \end{aligned}$$

If $|\psi\rangle$ is an e. vec of \vec{L}^2 with e-val $\hbar^2 l(l+1)$ (and of L_z with e-val $m\hbar$), then

$$\vec{L}^2 |\psi\rangle = \hbar^2 l(l+1) |\psi\rangle \quad \text{and} \quad L_z |\psi\rangle = m\hbar |\psi\rangle \quad (19)$$

and

$$\langle \vec{r} | \vec{p}^2 | \psi \rangle = -\frac{\hbar^2}{r^2} \partial_r (r^2 \partial_r) \langle \vec{r} | \psi \rangle + \frac{\hbar^2}{r^2} l(l+1) \langle \vec{r} | \psi \rangle. \quad (20)$$

This is the starting point for a discussion of central potentials $V = V(r)$.